

# Neutral Leptons in Cosmology

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BEYOND 2010, Cape Town

# Neutral Leptons in Cosmology or New Physics Without New Energy Scale

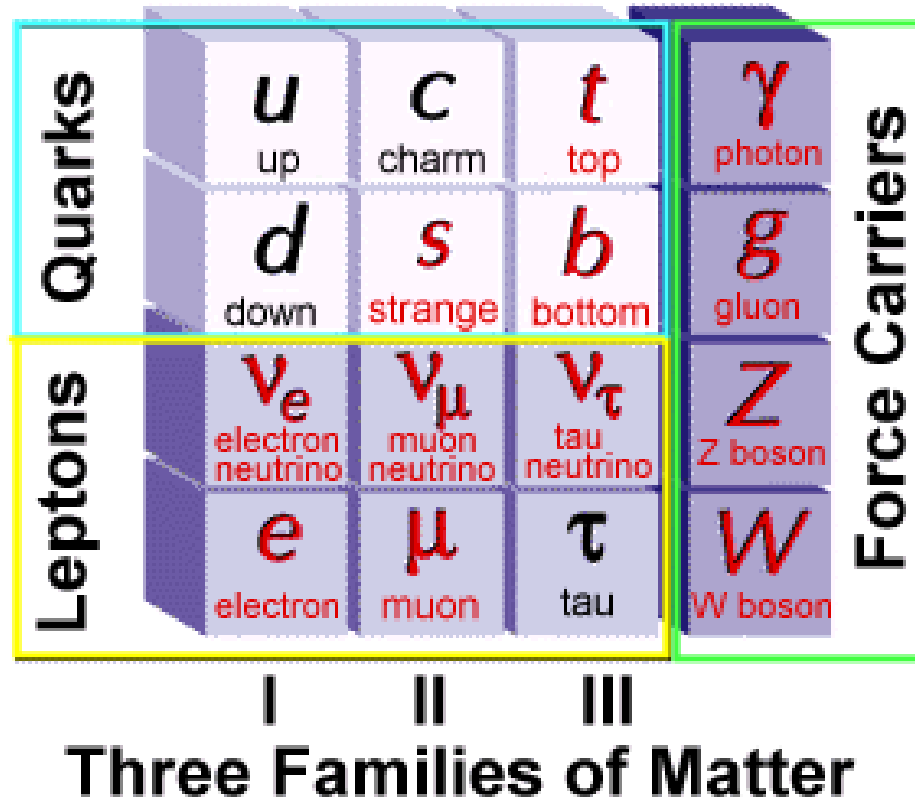
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- Motivation
- Neutrino masses
- Dark Matter
- Baryon Asymmetry
- Rare decays
- Conclusions

Standard Model of particle interactions is in great shape: it agrees with all **accelerator** experiments

# Elementary Particles



The only missing particle - the Higgs boson. It will be searched at the LHC

Still, the Standard Model cannot accommodate a number of cosmological observations and discoveries in neutrino physics, it also has a number of “fine tuning” problems from theory side.

## Our strategy

- Select the most important problems to solve (may be subjective).
- Use Ockham’s razor principle: “*Frustra fit per plura quod potest fieri per pauciora*” or “entities must not be multiplied beyond necessity”. For particle physics: entities = new hypothetical particles and new scales different from Fermi and Planck scales.

# SM problems and possible solutions

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Hierarchy problem: stability of the Higgs mass against radiative corrections

Possible solutions:

- Compensation of divergent diagrams by new particles at TeV scale (supersymmetry, composite Higgs boson). Consequence: new physics at LHC

## SM problems and possible solutions

- New symmetry – exact, but spontaneously broken scale invariance. Higgs mass is kept small in the same way as photon mass is kept zero by gauge invariance. Consequences: validity of the SM all the way up to the Planck scale, nothing but the Higgs at LHC in the mass interval

$$m_{\min} < m_H < m_{\max}$$

$$m_{\min} = \left[ 126.3 + \frac{m_t - 171.2}{2.1} \times 4.1 - \frac{\alpha_s - 0.1176}{0.002} \times 1.5 \right] \text{ GeV}$$

$$m_{\max} = \left[ 173.5 + \frac{m_t - 171.2}{2.1} \times 1.1 - \frac{\alpha_s - 0.1176}{0.002} \times 0.3 \right] \text{ GeV}$$

theory error in  $m_{\min} \simeq \pm 2$  GeV. Existence of massless particle – dilaton, which can play the role of **Dark Energy**

# SM problems and possible solutions

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The universe is flat, homogeneous and isotropic with high accuracy. It contained in the past small density fluctuations that lead to structure formation

Possible solution: inflation.

The inflaton (scalar particle inflating the universe) is

- new particle with the mass of the order of  $10^{13}$  GeV and minimal coupling to gravity

Alternative

- SM Higgs boson with non-minimal coupling to gravity



# SM problems and possible solutions

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## Neutrino masses and oscillations

Possible solutions:

- See-saw mechanism: existence of several superheavy ( $M \sim 10^{10}$  GeV) neutral leptons. Direct experimental consequences: none, as the mass is too large to be accessed

## Alternative

- Existence of new leptonic flavours with masses similar to those of known quarks and leptons. Experimental consequence: possibility of direct experimental search

# SM problems and possible solutions

## Dark matter

Possible solutions:

- WIMPS with masses of the order of 100 GeV and roughly electroweak cross-sections (e.g. SUSY neutralino).

Consequences: new particles at LHC, success of WIMP searches

## Alternative

- Super-WIMPS with masses in keV region. Natural possibility: new neutral leptonic flavour with mass of few keV. Consequences: no DM candidates at LHC, failure of WIMP searches. Possibility of search through radiative processes  $N \rightarrow \nu\gamma$  which leads to existence of narrow X-ray line in direction of DM concentrations.

# SM problems and possible solutions

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## Baryon asymmetry of the Universe

Possible solutions:

- Baryogenesis due to new physics above the electroweak scale. Potential consequences: new particles at LHC (for electroweak baryogenesis)

## Alternative

- Baryogenesis due to new neutral leptonic flavours with masses in the region from 140 MeV up to few GeV. Experimental consequence: possibility of direct experimental search

# Realisation: $\nu$ MSM

SM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	$\nu_e$	e	$\nu_\mu$	$\mu$	$\nu_\tau$	$\tau$
right		e		$\mu$		$\tau$

leptons

$\nu$ MSM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	$\nu_e$	e	$\nu_\mu$	$\mu$	$\nu_\tau$	$\tau$
right	$N_e$	e	$N_\mu$	$\mu$	$N_\tau$	$\tau$

leptons

**Role** of  $N_e$  with mass in keV region: dark matter

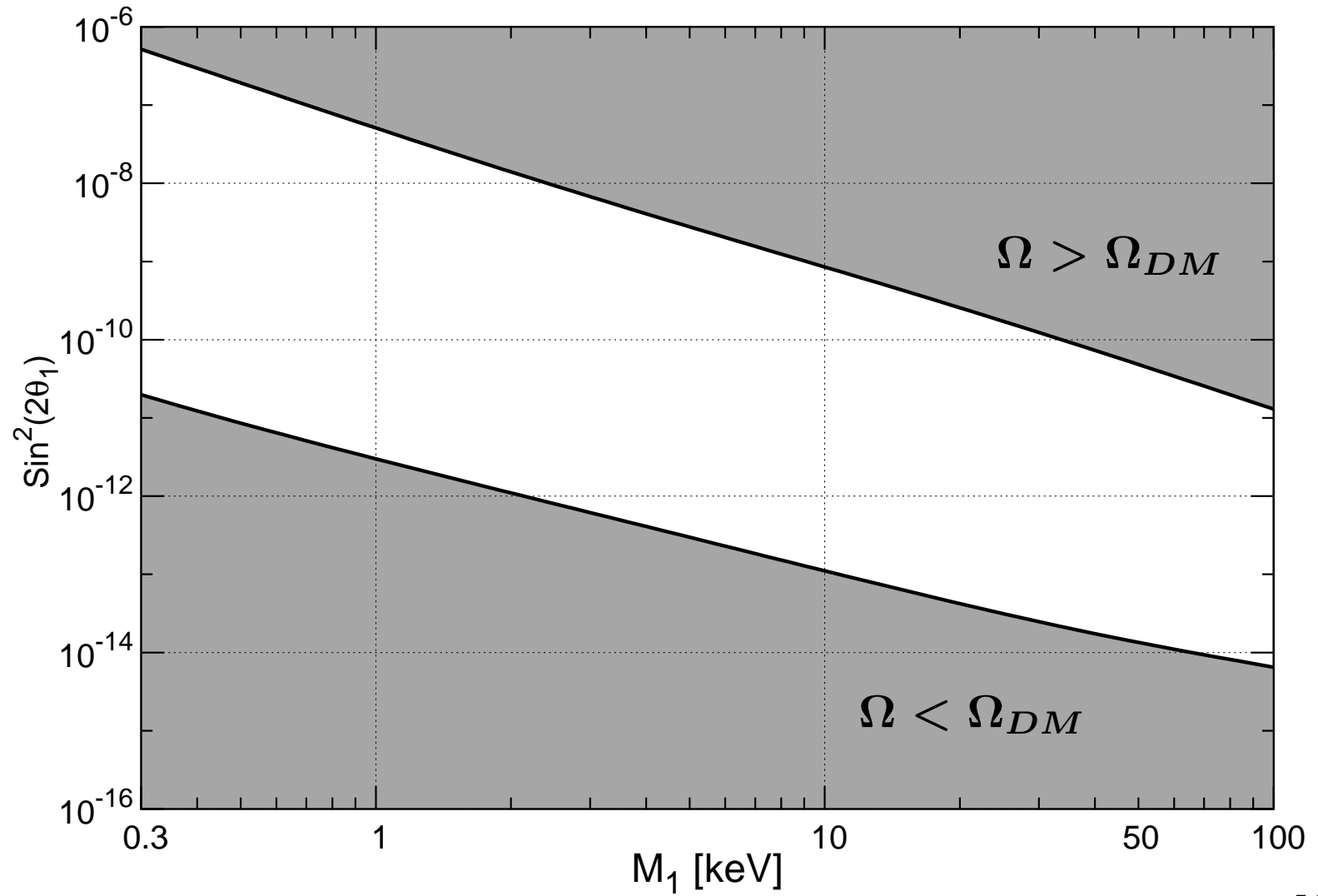
**Role** of  $N_\mu$ ,  $N_\tau$  with mass in 100 MeV – GeV region: “give” masses to neutrinos and produce baryon asymmetry of the Universe

**Role** of the Higgs: give masses to quarks, leptons,  $Z$  and  $W$  and inflate the Universe.

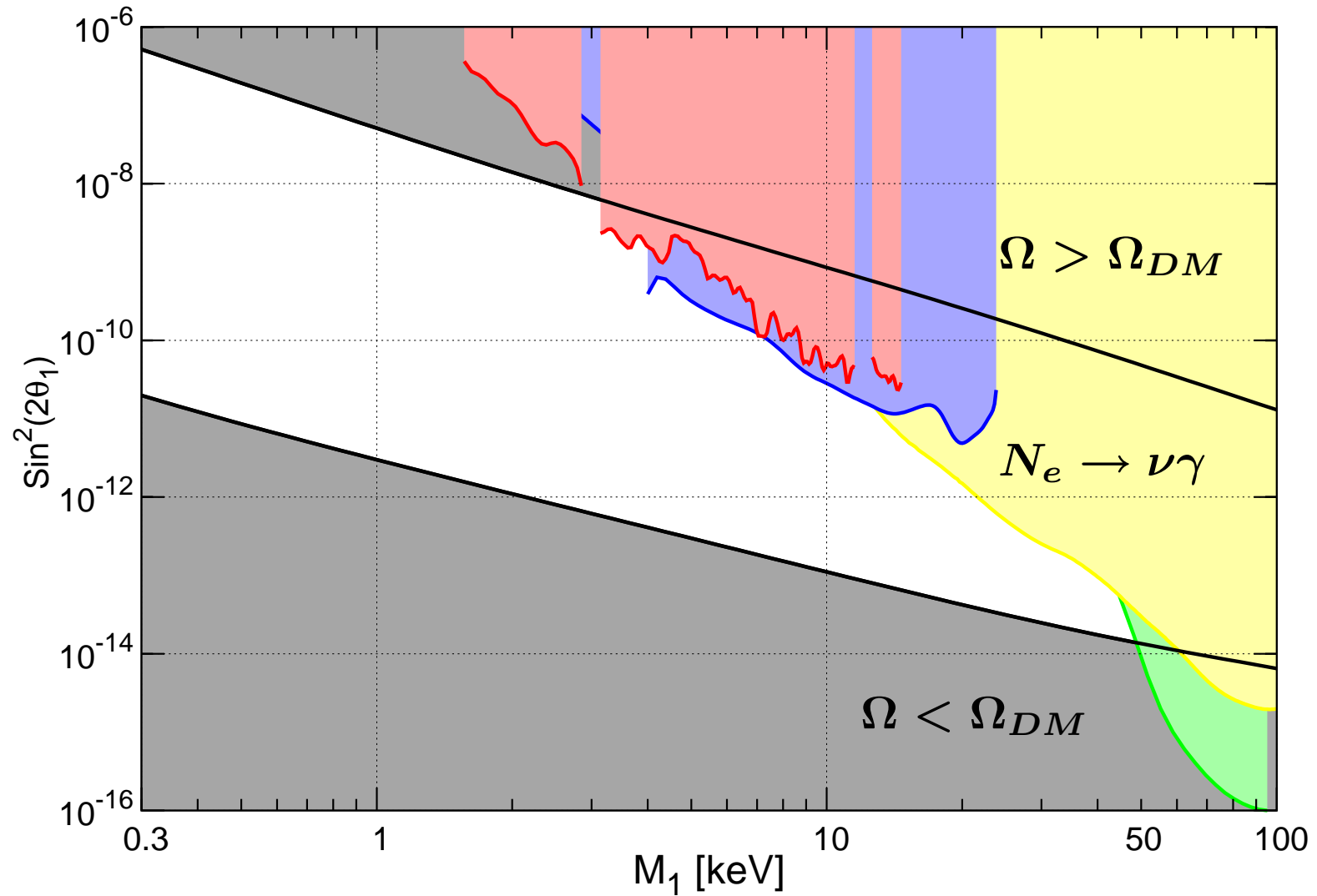
# Constraints on DM sterile neutrino

- **Production.**  $N_e$  are created in the early Universe in reactions  $l\bar{l} \rightarrow \nu N_e$ ,  $q\bar{q} \rightarrow \nu N_e$  etc. We should get correct DM abundance.
- **X-rays.**  $N_e$  decays radiatively,  $N_e \rightarrow \gamma\nu$ , producing a narrow line which can be detected. This line has not been seen (yet).
- **Structure formation.** If  $N_e$  is too light it may have considerable free streaming length and erase fluctuations on small scales. This can be checked by the study of Lyman- $\alpha$  forest spectra of distant quasars.

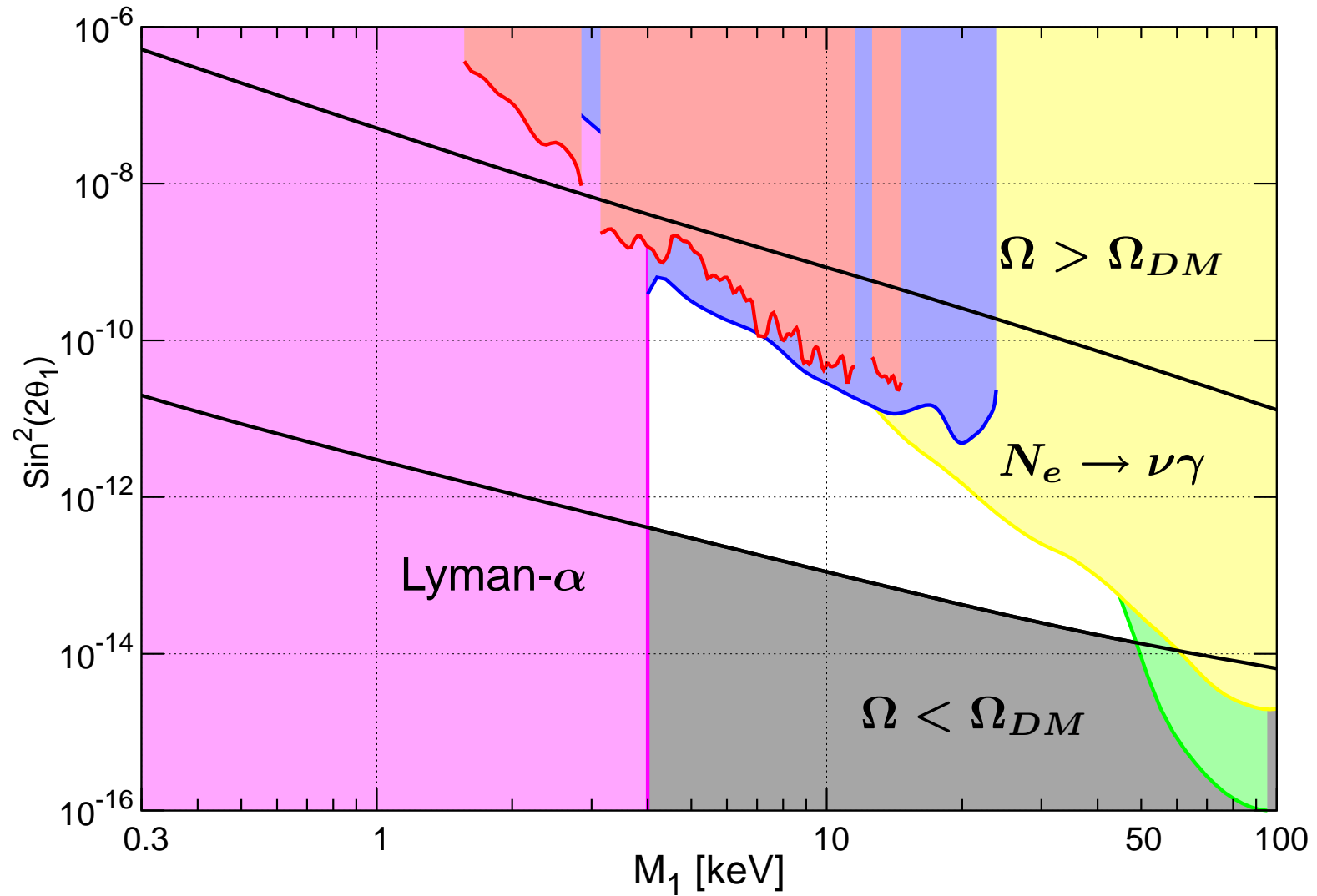
# DM: production



# DM: production + X-ray constraints



# DM: production + X-ray constraints + Lyman- $\alpha$ bounds





## Direct experimental test: astrophysics

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Over the last years restrictions on sterile neutrino parameters were improved **by several orders of magnitude**.

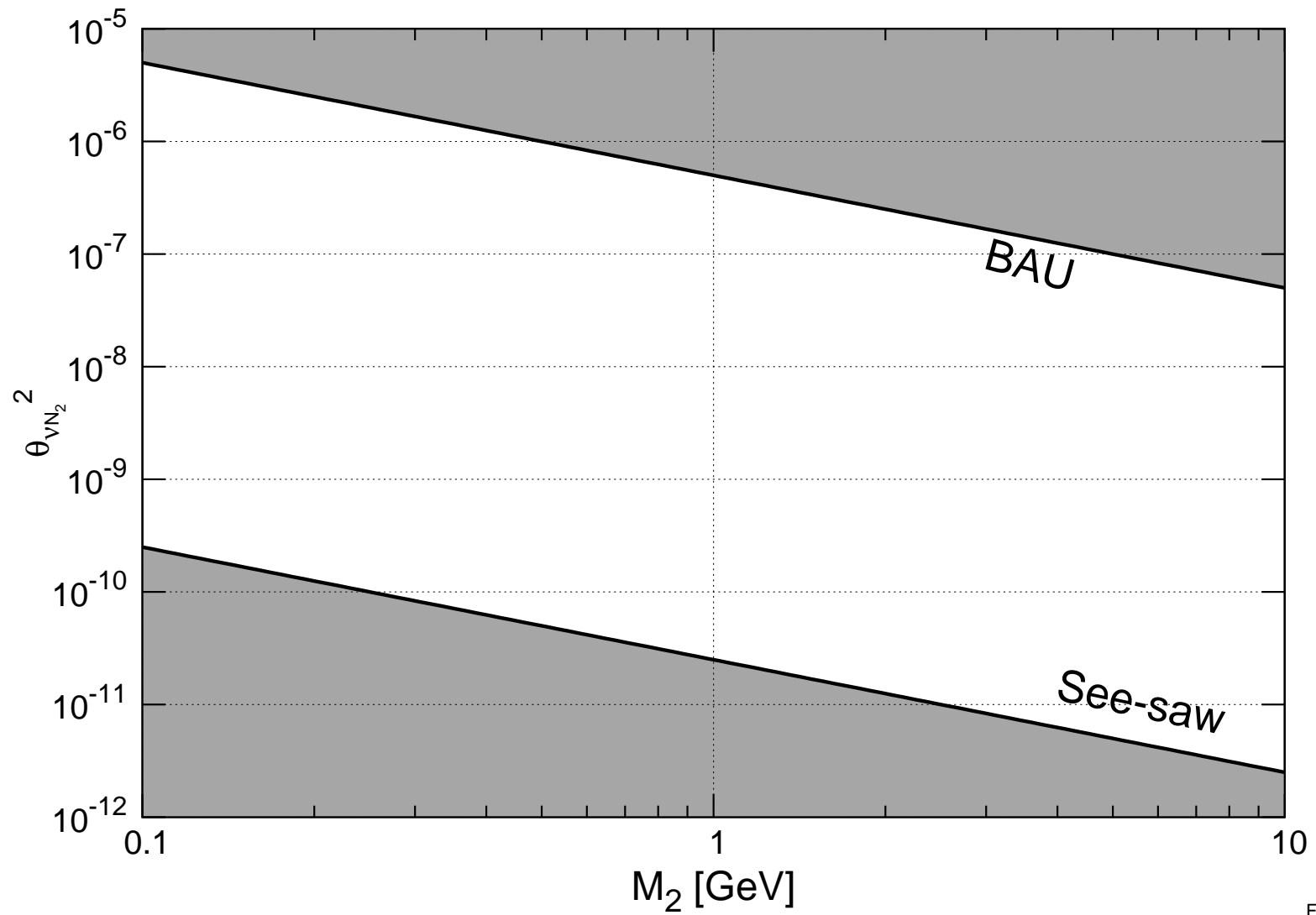
The new data from *Chandra* and *XMM-Newton* can hardly improve constraints by more than a factor 10. One needs:

- Improvement of spectral resolution up to the natural line width ( $\Delta E/E \sim 10^{-3}$ ).
- $\text{FoV} \sim 1^\circ$  (size of a dSph).
- Wide energy scan, from  $\mathcal{O}(100)$  eV to  $\mathcal{O}(10)$  MeV.

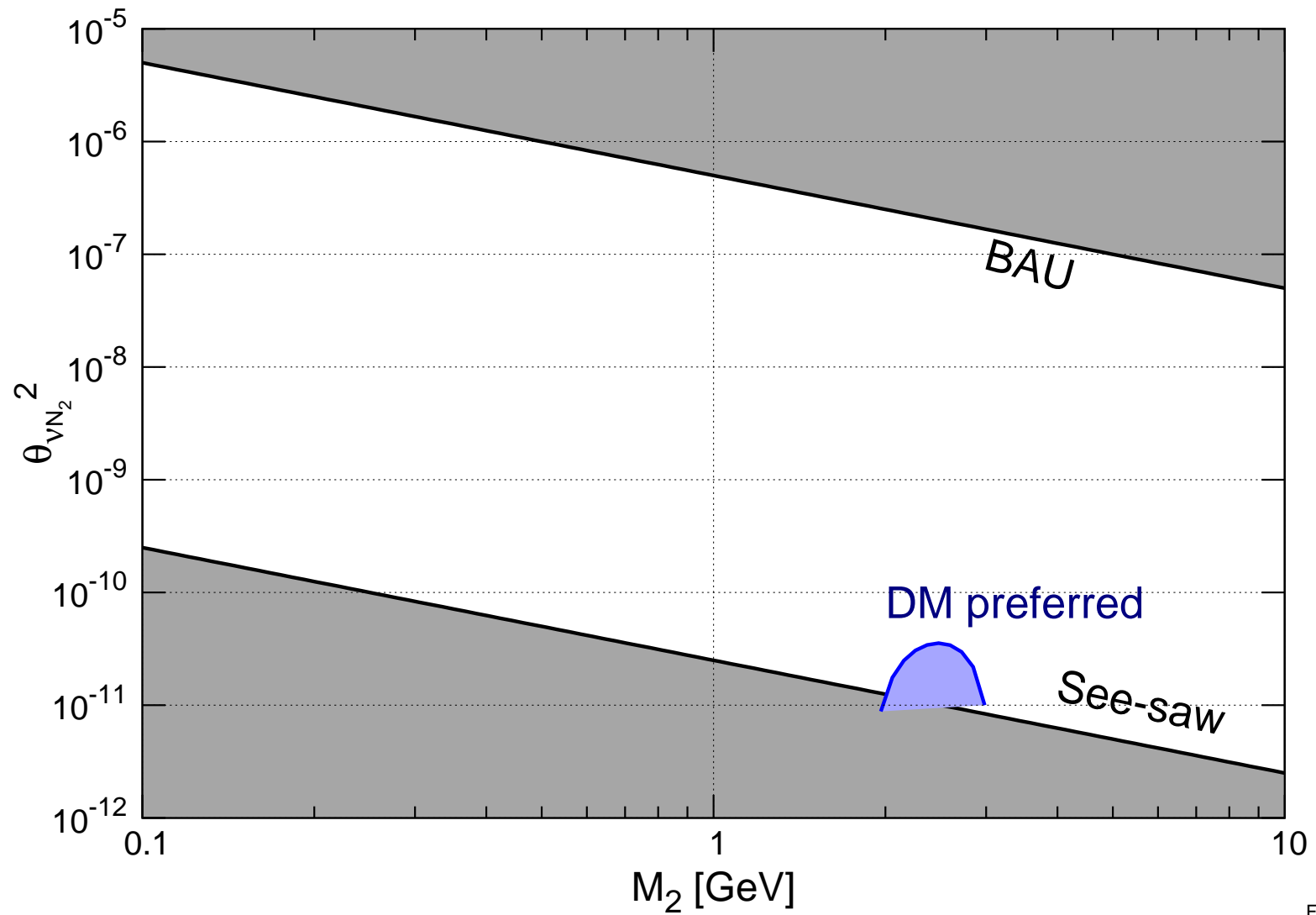
# Constraints on BAU Majorana fermions

- **BAU generation** requires out of equilibrium: mixing angle of  $N_{\mu,\tau}$  to active neutrinos cannot be too large
- **Neutrino masses.** Mixing angle of  $N_{\mu,\tau}$  to active neutrinos cannot be too small
- **Dark matter and BAU.** Concentration of DM sterile neutrinos must be much larger than concentration of baryons
- **BBN.** Decays of  $N_{\mu,\tau}$  must not spoil Big Bang Nucleosynthesis
- **Experiment.**  $N_{\mu,\tau}$  have not been seen (yet).

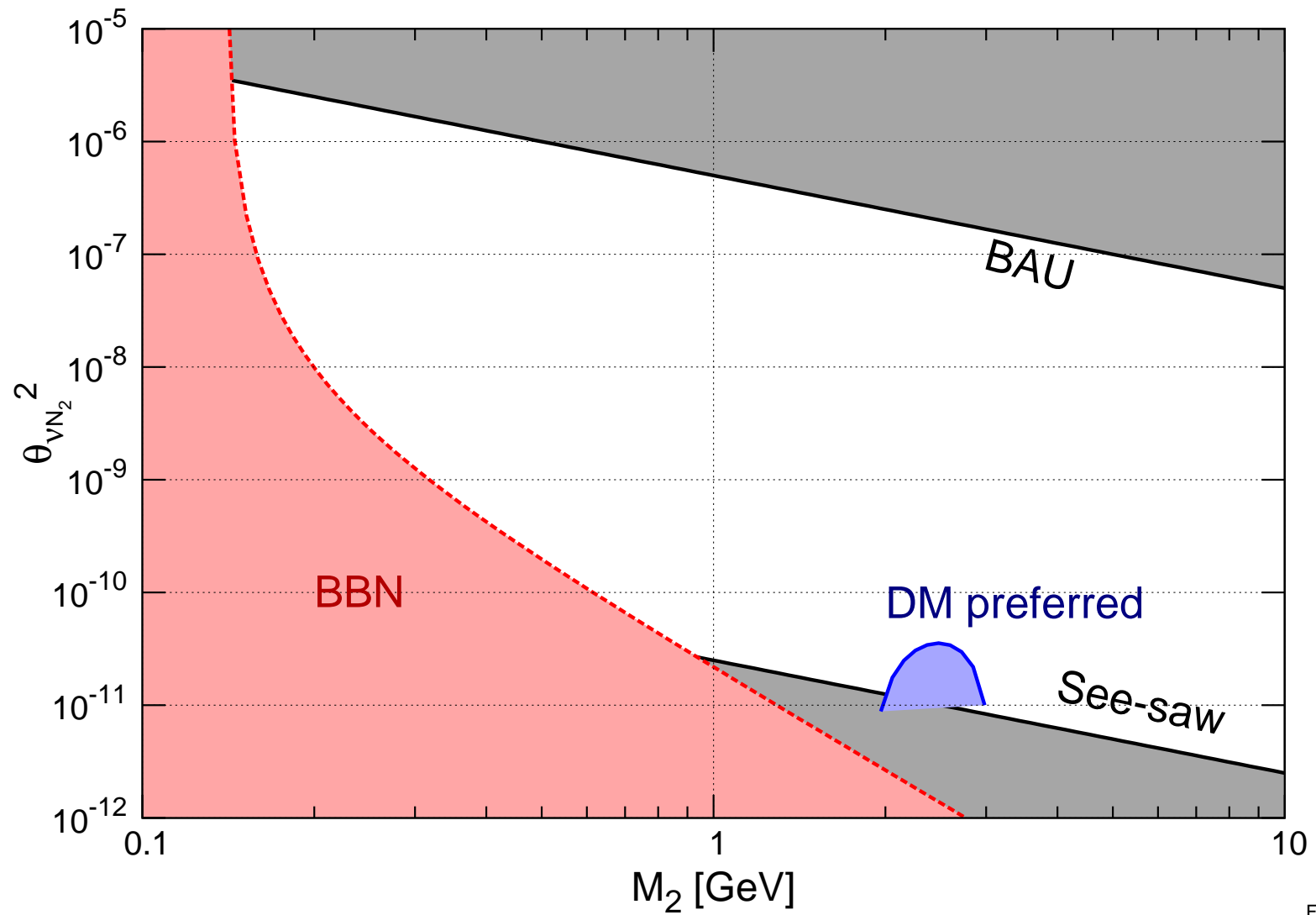
# $N_{\mu,\tau}$ : BAU



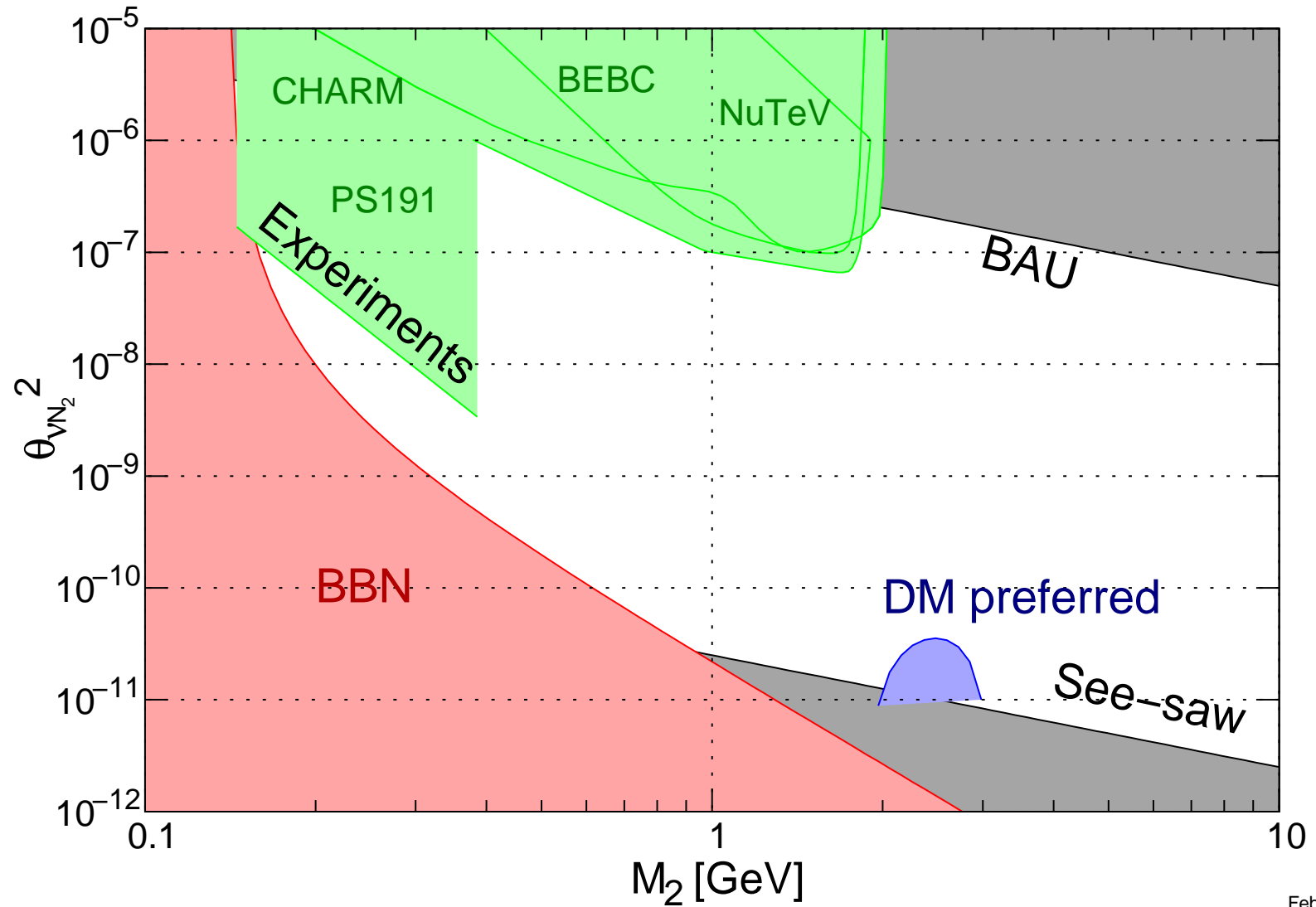
# $N_{\mu,\tau}$ : BAU + DM



# $N_{\mu,\tau}$ : BAU + DM + BBN



# $N_{\mu,\tau}$ : BAU + DM + BBN + Experiment



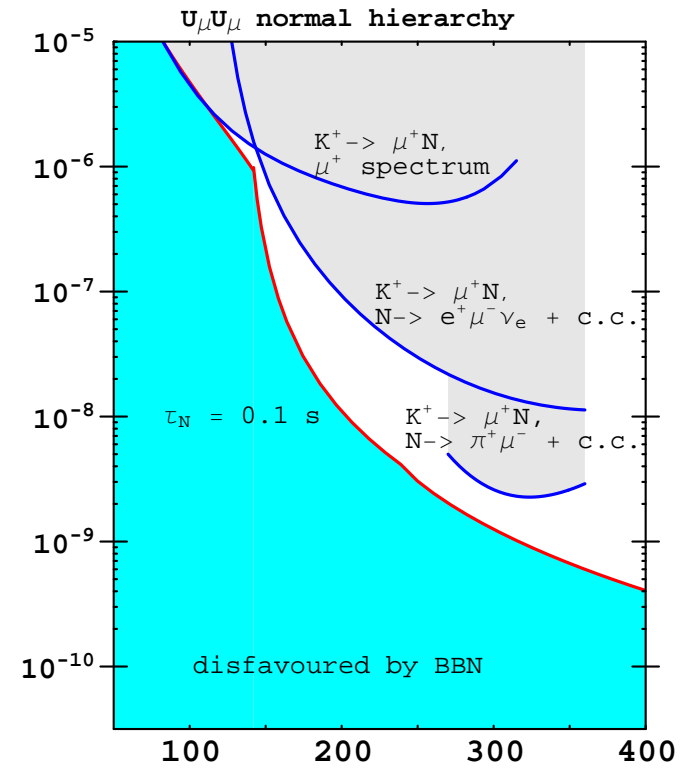
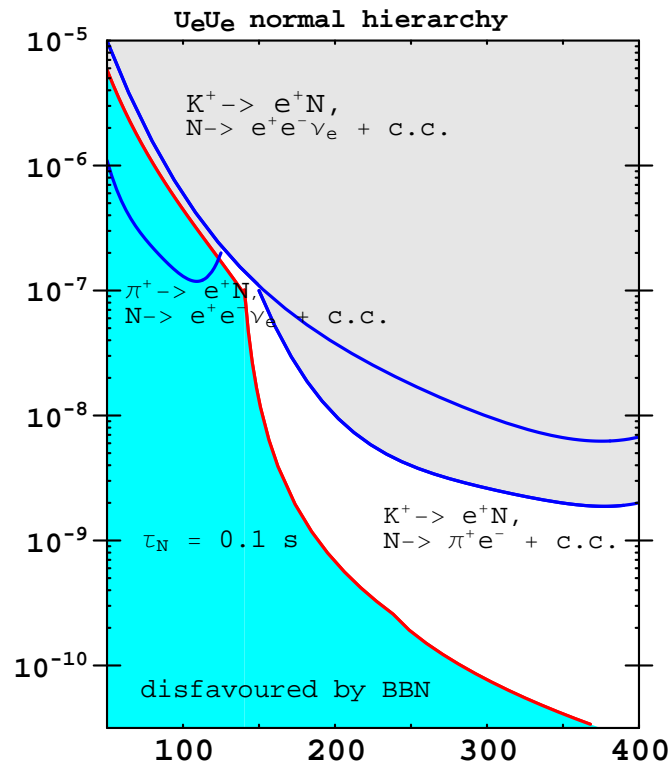
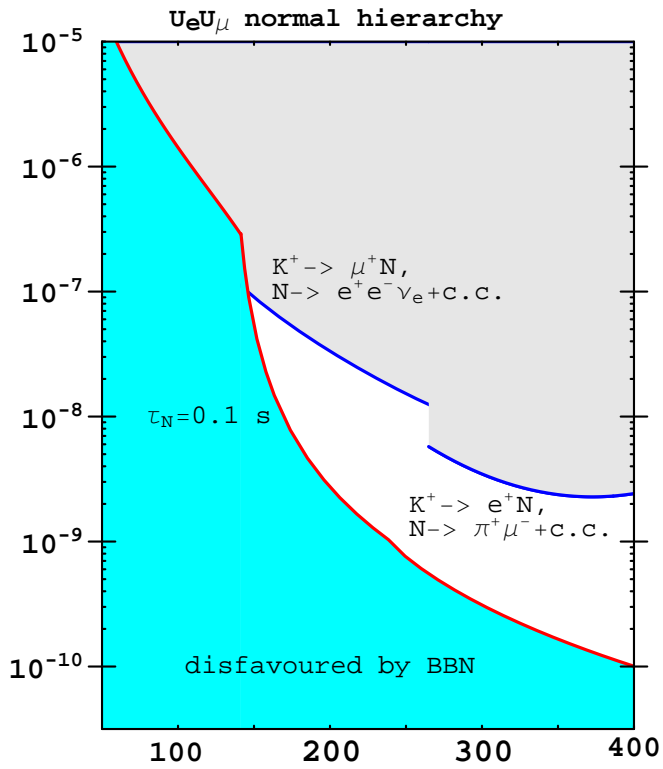
# Direct experimental tests: rare decays

# Previous searches at CERN

- A. M. Cooper-Sarkar *et al.* [WA66 Collaboration] “Search For Heavy Neutrino Decays In The Bebc Beam Dump Experiment”, 1985
- J. Dorenbosch *et al.* [CHARM Collaboration] “A search for decays of heavy neutrinos in the mass range 0.5-GeV to 2.8-GeV”, 1985
- G. Bernardi *et al.* [PS191 Collaboration], “Search For Neutrino Decay”, 1986;  
“Further Limits On Heavy Neutrino Couplings”, 1988
- P. Astier *et al.* [NOMAD Collaboration], “Search for heavy neutrinos mixing with tau neutrinos”, 2001
- P. Achard *et al.* [L3 Collaboration], “Search for heavy neutral and charged leptons in  $e^+e^-$  annihilation at LEP”, 2001



# CERN PS191 experiment, 1988



Conclusion:  $M_{2,3} > 140 \text{ MeV}$

# Experimental signatures 1

Challenge - from baryon asymmetry:  $\theta^2 \lesssim 5 \times 10^{-7} \left(\frac{\text{GeV}}{M}\right)$

- Peak from 2-body decay and missing energy signal from 3-body decays of  $K$ ,  $D$  and  $B$  mesons (sensitivity  $\theta^2$ )

Example:

$$K^+ \rightarrow \mu^+ N, \quad M_N^2 = (p_K - p_\mu)^2 \neq 0$$

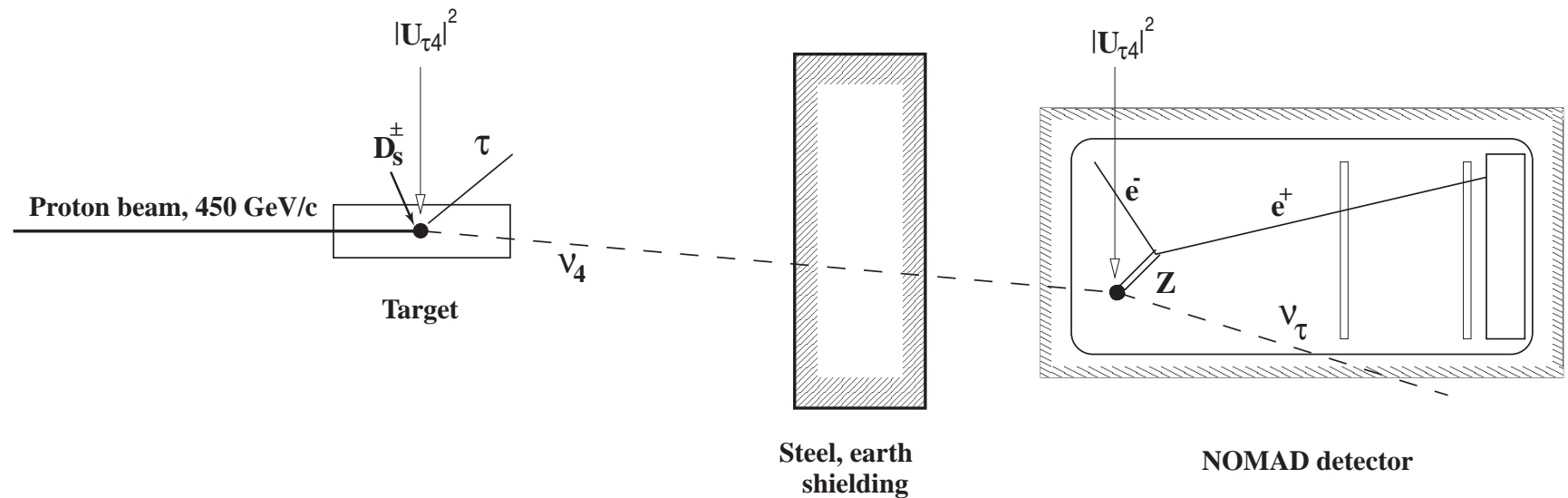
Similar for charm and beauty.

- $M_N < M_K$ : KLOE, NA62, E787
- $M_K < M_N < M_D$ : charm and  $\tau$  factories, CLEO
- $M_N < M_B$ : B-factories (planned luminosity is not enough to get into cosmologically interesting region)

# Experimental signatures 2

- Two charged tracks from a common vertex, decay processes  $N \rightarrow \mu^+ \mu^- \nu$ , etc. (sensitivity  $\theta^4 = \theta^2 \times \theta^2$ )  
**First step:** proton beam dump, creation of  $N$  in decays of  $K$ ,  $D$  or  $B$  mesons:  $\theta^2$   
**Second step:** search for decays of  $N$  in a near detector, to collect all  $N$ s:  $\theta^2$ 
  - $M_N < M_K$ : Any intense source of K-mesons (e.g. from proton targets of PS or SPS.)
  - $M_N < M_D$ : SPS beam + near detector
  - $M_N < M_B$ : Very high intensity p-accelerator ( $E \sim 1\text{TeV}$ ) + near detector
  - $M_N > M_B$ : extremely difficult

# $N_{2,3}$ production and decays



Type on neutrino mass hierarchy - from branching ratios of  $N_{2,3}$  decays to  $e, \mu, \tau$ .

CP asymmetry can be as large as 1% - from BAU and DM

# Conclusions

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- New physics, responsible for **neutrino masses and mixings**, for **dark matter**, and for **baryon asymmetry of the Universe** may hide itself **below** the EW scale

# Conclusions

- New physics, responsible for **neutrino masses and mixings**, for **dark matter**, and for **baryon asymmetry of the Universe** may hide itself **below** the EW scale
- Search for new physics: very high intensity proton accelerators, very high intensity B or charm factories, rather than high energy colliders. SPS at CERN is a very good opportunity. Prediction: no new physics at LHC (except Higgs), ILC or muon collider.

